# ALMOST MARKOV PROCESSES IN MONTE CARLO
# SIMULATION OF BIOLOGICAL MOLECULES

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# Almost Markov Processes in Monte Carlo Simulation of Biological Molecules

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**Abstract**

Computer simulation is an important tool for investigating prop-
erties of biological molecules. This paper presents newly-developed
optimization techniques that take into account the inhomogeneity
and anisotropy essential to the structure and function of biological
molecules to produce highly efficient simulations.

During the past few decades, computer simulations have become an in-
creasingly important tool for the study of thermodynamic systems [1]. Ex-
tremely rapid advances in both hardware and simulation techniques have
improved the accuracy of such calculations and broadened the range of prob-
lems considered.

An important area of computer simulation at the borders of medicine,
biology, chemistry, and physics is the study of macromolecules [2, 3]. Because
of the large number of atoms and the complexity of the interactions, even
today's largest supercomputers are not yet sufficiently powerful to deal with
all the problems of interest. Consequently, new approaches to these problems
are needed to increase the efficiency of these calculations.

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In this paper, we present a new approach we have developed. We have achieved substantial improvement in both the efficiency and reliability of the simulations, but our approach has also opened up some interesting mathematical questions about the nature of an approximation we have introduced.

The computer simulation methods currently in use (molecular dynamics and Monte Carlo) were originally developed for treating fluids. They naturally made use of the simplifying feature that fluids are isotropic and homogeneous. Unfortunately, macromolecules have neither property, which leads to a large decrease in efficiency.

A Monte Carlo simulation stores an explicit configuration of the system, in this case, a large molecule. The algorithm is a stochastic process that consists of moving atoms, or groups of atoms, to simulate the thermal motion of the system. It is usual to consider only Markov processes, in which the transition probabilities depend only on the current configuration of the system, and not on its history. This enables the use of a theorem that if a condition on the transition probabilities is satisfied (detailed balance) and the system is ergodic (any configuration can be reached from any other configuration in a finite number of steps with non-zero probability), then the simulation will correctly reproduce the thermal distribution (in the limit of infinite computer time).

In setting up a simulation, there are generally one or two global parameters that must be specified to optimize the efficiency. These are normally determined during a preliminary simulation, and then fixed during the simulation proper.

In order to take into account the inhomogeneity and the local anisotropy, we have used data accumulated during the simulation to determine parameters for optimizing moves of individual atoms. This necessarily relaxes the Markov condition, because the new transition probabilities depend on the recent history of the simulation [4]. Empirical studies demonstrate that the almost Markovian nature of the simulation does not introduce any systematic errors, but a mathematical proof is still missing.

To develop the new method, we approximate the local potential by an anisotropic simple harmonic oscillator of the form

$$H = \frac{1}{2} \sum_{ij} k_{ij} x_i x_j \qquad (1)$$

where $k_{ij}$ represents the spring constants. A transformation matrix, $D_{ij}$, is

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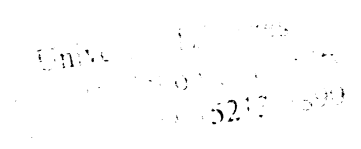
used to generate moves in an ellipsoid given by

$$d_i = \sum_j D_{ij} v_j \qquad (2)$$

where $v_j$ is a random vector chosen from a unit sphere. The optimum jump size scales as the contours of constant energy and is determined by a dimensionless parameter, $F$, given by

$$F^2 = \frac{1}{2}\beta \sum_{ij} k_{ij} d_i d_j \qquad (3)$$

where $\beta = \frac{1}{k_B T}$. $F$ is a scale factor chosen to optimize the efficiency of the simulation. The effective values of $k_{ij}$ are determined from the simulation using

$$[\Delta E d_l d_m] = \frac{1}{2}\sum_{ij} k_{ij}[d_i d_j d_l d_m] \qquad (4)$$

where $\Delta E$ denotes the change of energy for an attempted move $d_i$ and the square brackets indicate an average over all attempted moves, whether or not they were accepted. The matrix $D_{in}$ is then obtained from

$$D_{in} = F\sqrt{\frac{2}{\beta \lambda_n}} V_{in} \qquad (5)$$

where $\lambda_n$ is an eigenvalue of $K_{ij}$ and $V_{in}$ is the corresponding normalized eigenvector. $D_{ij}$ is updated periodically to adapt to the changing local environment of each atom. This method, which we denote as Dynamically Optimized Monte Carlo ($DOMC$) turns out to be remarkably robust when applied to a wide variety of Hamiltonians.

As a simple test for systematic error, we simulated a one-dimensional simple harmonic oscillator, and used DOMC with very short cycles. Averaging over only 1 MC-step/cycle, we find a large systematic error of 42%. However, with even 2 MC-steps/cycle, the error is reduced to 3%, and for 3 MC-steps/cycle, it is just over 1%. No systematic error was measurable for 4 MC-steps/cycle. Since we have generally used more than 10 MC-steps/cycle in our work, the systematic error appears to be negligible.

We have already applied this method to a large number of systems, including highly anisotropic oscillators and biomolecules ranging from adenosine (one of the bases of $DNA$) to complex proteins such as the Bovine

Pancreatic Trypsin Inhibitor (*BPTI*). Our programs have included additional refinements beyond the scope of this paper to take long-wave-length fluctuations into account more effectively. Results have been very good, with large improvements in efficiency.

Work is continuing to extend this approach and prospects are extremely promising. Efforts to analyze the almost Markovian nature of these algorithms and provide a clear understanding and estimation of the systemmatic errors is an important aspect of this research.

# References

[1] *Applications of the Monte Carlo Method in Statistical Physics.* Vol. 36, 2nd ed., edited by K. Binder (Springer, Berlin, 1987).

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[4] S. Shumway and J. P. Sethna (Bull. Am. Phys. Soc. **35** 500 (1990)) have presented a novel technique for optimizing MC moves. It is different from ours, but it also uses an almost Markovian process to optimize the simulations.

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