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## NAMT <br> $91-001$

# Jacobians and Hardy Spaces 

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Research Report No. 91-NA-001
May 1991

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## Introduction

We wish to illustrate here the links between various nonlinear quantities identified by the compensated-compactness method (L. Tartar [9],[10]), F. Murat [6],[7]) and the Hardy spaces. A systematic study of this matter can be found in R. Coifman, P. L. Lions, Y. Meyer and S. Semmes [3], and we want to illustrate it on the particular example of the Jacobian.

More precisely, we consider $J(u)=\operatorname{det}(\nabla u)$ when $u \in W^{1, p}\left(\mathbb{R}^{N}\right)^{N}$ for some $p \in[1 . \infty[$ where $N \geq 2$. This quantity clearly makes sense in $L^{1}\left(\mathbb{R}^{N}\right)$ if $p=N$; however, in that case, there are various reasons to guess that $J(u)$ might be slightly better than $L^{1}$. First of all, it is a classical fact that $J$ is weakly sequentially continuous on $W^{1, N}\left(\mathbb{R}^{N}\right) N$ - this fact is one of the key ingredients in J. Ball's theory of polyconvex functionals in Nonlinear Elasticity [1]. Next, several results by H. Wente [12], L. Tartar [11] indicate that $\mathrm{L}^{1}$ is not optimal - see also H . Brézis and J. M. Coron [2]. Finally, the last piece of evidence is a striking recent result due to S. Müller [4] showing that if $\left.u \in W_{10 c}^{1, N_{(\mathbb{R}}}{ }^{N}\right) N$ and $J(u) \geq 0$ a.e., then $J(u) \log (1+J(u)) \in$ $L_{\text {loc }}^{1}\left(\mathbb{R}^{N}\right)$.

We will show below in Section 2 how all these results can be recovered from the following statement: $J(u) \in \mathscr{H}^{1}\left(\mathbb{R}^{N}\right)$ if $u \in W^{1, N}\left(\mathbb{R}^{N}\right)^{N}$. Here and everywhere below, we denote by $\mathscr{O}^{\mathrm{p}}\left(\mathbb{R}^{\mathrm{N}}\right)$ the following Hardy spaces:

$$
\begin{equation*}
\mathscr{H}^{\mathrm{p}}\left(\mathbb{R}^{\mathrm{N}}\right)=\left\{\mathrm{f} \in \mathrm{~S}^{\prime}\left(\mathbb{R}^{\mathrm{N}}\right) / \sup _{\mathrm{t}>0}\left|\mathrm{~h}_{\mathrm{t}} * \mathrm{f}\right| \in \mathrm{L}^{\mathrm{p}}\left(\mathbb{R}^{\mathbf{N}}\right)\right\}, \mathrm{p}>0, \tag{1}
\end{equation*}
$$

where $h_{t}$ is a regularization kernel satisfying, for example

$$
\begin{equation*}
\mathrm{h}_{\mathrm{t}}=\frac{1}{\mathrm{t}^{\mathbf{N}}} \mathrm{h}(\dot{\bar{t}}), \mathrm{h} \in \mathscr{D}\left(\mathbb{R}^{\mathrm{N}}\right), 0 \leq \mathrm{h} \text { on } \mathbb{R}^{\mathrm{N}}, \text { Supp } \mathrm{h} \subset \mathrm{~B}_{1}, \tag{2}
\end{equation*}
$$

where we denote by $B_{\lambda}$ the open ball of radius $\lambda$ and by $B_{\lambda}(x)=B(x, \lambda)$ the open ball centered at x of radius $\lambda\left(\mathrm{B}_{\lambda}=\mathrm{B}_{\lambda}(0)=\mathrm{B}(0, \lambda)\right)$. We shall see in fact that $\mathrm{J}(\mathrm{u}) \in \mathscr{H}^{\mathrm{p}}\left(\mathbb{R}^{\mathrm{N}}\right)$ if $\mathrm{u} \in \mathrm{W}^{1, \mathrm{pN}}\left(\mathbb{R}^{\mathrm{N}}\right)^{\mathrm{N}}, \frac{\mathrm{N}}{\mathrm{N}+1}<\mathrm{p} \leq 1$ and that $\mathrm{J}(\nabla \phi)=\operatorname{det}\left(\mathrm{D}^{2} \phi\right) \in \mathscr{H}^{\mathrm{p}}\left(\mathbb{R}^{\mathrm{N}}\right)$ if $\phi$ $\epsilon W^{2, p N}\left(\mathbb{R}^{N}\right)^{N}, \frac{N}{N+2}<p \leq 1$. Of course, one has to interpret $J(u)$ or $J(\nabla \phi)$ in a distribution sense explained below when $\mathrm{p}<1$.

We present in Section 3 a proof of these facts. To simplify the presentation and the notations we restrict our attention here to the case when $\mathrm{N}=2$.

## 2. Main results and consequences.

Of course, if $u \in W^{1,2}\left(\mathbb{R}^{2}\right)^{2}, J(u)=\operatorname{det}(\nabla u) \in L^{2}\left(\mathbb{R}^{2}\right)$. But, one easily checks that

$$
\begin{equation*}
\mathrm{J}(\mathrm{u})=\partial_{1}\left(\mathrm{u}_{1} \partial_{2} \mathrm{u}_{2}\right)-\partial_{2}\left(\mathrm{u}_{2} \partial_{1} \mathrm{u}_{2}\right) \text { in } \mathscr{D}^{\prime}\left(\mathbb{R}^{2}\right) \tag{3}
\end{equation*}
$$

and this last expression is well-defined (in the sense of distributions) whenever $u \in W^{1, \frac{4}{3}}\left(\mathbb{R}^{2}\right)^{2}$ : indeed, one then deduces from Sobolev embeddings that $u \in L^{4}\left(\mathbb{R}^{2}\right)^{2}$ and thus $|u||\nabla u|$ is integrable.

Next, if $u=\nabla \phi$, further cancellations of $J(u)$ take place and we may write

$$
\begin{equation*}
\mathrm{J}(\nabla \phi)=-\frac{1}{2} \partial_{1}^{2}\left(\left(\partial_{2} \phi\right)^{2}\right)-\frac{1}{2} \partial_{2}^{2}\left(\left(\partial_{1} \phi\right)^{2}\right)+\partial_{12}^{2}\left(\partial_{1} \phi \partial_{2} \phi\right) \text { in } \mathscr{D}^{\prime}\left(\mathbb{R}^{2}\right) \tag{4}
\end{equation*}
$$

or

$$
\begin{equation*}
\mathrm{J}(\nabla \phi)=\frac{1}{2} \partial_{1}^{2}\left(\phi \partial_{2}^{2} \phi\right)+\frac{1}{2} \partial_{2}^{2}\left(\phi \partial_{1}^{2} \phi\right)-\partial_{12}^{2}\left(\phi \partial_{12}^{2} \phi\right) \text { in } \mathscr{D}^{\prime}\left(\mathbb{R}^{2}\right) . \tag{5}
\end{equation*}
$$

Note that (4) makes sense as soon as $\phi \in \mathrm{W}^{1,2}\left(\mathbb{R}^{2}\right)$ while (5) makes sense if $\phi \in \mathrm{W}^{2,1}\left(\mathbb{R}^{2}\right)$ (since it implies $\phi \in \mathrm{C}_{0}\left(\mathbb{R}^{2}\right)$ ). In fact, the two expressions are easily shown to be equal if $\phi \in \mathrm{W}^{2,1}\left(\mathbb{R}^{2}\right)$ since $\mathrm{W}^{2,1}\left(\mathbb{R}^{2}\right)$ embeds into $\mathrm{W}^{1,2}\left(\mathbb{R}^{2}\right)$ and $\partial_{\mathrm{j}}\left(\phi \partial_{\mathrm{j}} \phi\right)=\phi \partial_{\mathrm{j}}^{2} \phi+\left(\partial_{\mathrm{j}} \phi\right)^{2}\left(\forall_{\mathrm{j}}\right)$, $\partial_{1} \phi \partial_{2} \phi+\phi \partial_{12}^{2} \phi=\frac{1}{2} \partial_{1}\left(\phi \partial_{2} \phi\right)+\frac{1}{2} \partial_{2}\left(\phi \partial_{1} \phi\right)$. Notice, finally, that by Sobolev embeddings, (4) makes sense if $\phi \in \bar{W}^{2,1}\left(\mathbb{R}^{2}\right)=\left\{\phi \in W^{1,1}\left(\mathbb{R}^{2}\right) / \partial_{i, j}^{2} \phi\right.$ is a bounded measure on $\left.\mathbb{R}^{2}\left(\forall_{i, j}\right)\right\}$.

We may now state our main results:

Theorem 1: Let $\mathrm{p} \in\left(\frac{4}{3}, 2\right]$ and let $u \in \mathrm{~W}^{1, \mathrm{p}}\left(\mathbb{R}^{2}\right)^{2}$, then $\mathrm{J}(\mathrm{u}) \in \mathscr{\mathscr { O }}^{\mathrm{p} / 2}\left(\mathbb{R}^{2}\right)$.

Theorem 2: Let $\mathrm{p} \in(1,2]$ and let $\phi \in \mathrm{W}^{2, \mathrm{p}}\left(\mathbb{R}^{2}\right)$, then $\mathrm{J}(\nabla \phi) \in \mathscr{\mathscr { H }}^{\mathrm{P} / 2}\left(\mathbb{R}^{2}\right)$.

Remarks: 1) These results also hold locally.
2) The borderline cases $\mathrm{p}=\frac{4}{3}$ or $\mathrm{p}=1$ can also be studied - see also section 3 below.
3) One recovers immediately $S$. Müller's result [4] from Theorem 1 since if $f \in \mathrm{~L}_{\text {loc }}^{1}, f \geq 0$ a.e. then $f \in \mathscr{O}_{1 \text { oc }}^{1}\left(\mathbb{R}^{N}\right)$ if and only if $f \log (1+f) \in \mathrm{L}_{\text {loc }}^{1}$ (see E. Stein [8]).
4) Observing that $\mathscr{F}^{1}\left(\mathbb{R}^{2}\right) \in \mathrm{W}^{-1,2}\left(\mathbb{R}^{2}\right)$, one deduces that $W=(-\Delta)^{-1} J(u) \in W^{1,2}\left(\mathbb{R}^{2}\right)$ and $J(u) \in W^{-1,2}\left(\mathbb{R}^{2}\right)$ if $u \in W^{1,2}\left(\mathbb{R}^{2}\right)^{2}$. Furthermore, in that case, one sees that $\partial_{i, j} W \in \mathscr{F}^{1}\left(\mathbb{R}^{2}\right)\left(\forall_{i, j}\right)$ and this yields: $W \in \mathscr{F} L^{1}\left(\mathbb{R}^{2}\right)$. We recover in this way the results mentioned in the Introduction.
5) One can define a linear continuous map $P$ from $\mathscr{E}^{q}\left(\mathbb{R}^{N}\right)$ into $L^{q}\left(\mathbb{R}^{N}\right)$ for $0<q<1$ which consists in taking the "a. e. part" of a distribution $f$ in $\mathscr{F}^{q}\left(\mathbb{R}^{N}\right)$ : more precisely, $\operatorname{Pf}=\lim _{t \rightarrow 0}$ a.e. $h_{t} * f$. Of course, $\operatorname{Pf}=f$ if $f \in \mathscr{F}^{1}$ or $L^{1}$ and $\operatorname{Pf}$ is the regular part of $f$ if $f$ is a bounded measure.
Next, one remarks that when $\mathrm{p}<2$, one can also define a.e. $\operatorname{det}(\mathrm{zu})$ or $\operatorname{det}\left(\mathrm{D}^{2} \phi\right)$ obtaining thus a measurable function which lies obviously in $\mathrm{L}^{\mathrm{p} / 2}$. We denote by $\operatorname{Det}(\mathrm{\nabla u})$ or $\operatorname{Det}\left(\mathrm{D}^{2} \phi\right)$ these functions. Then, the above results yield easily: $\operatorname{Det}(\nabla \mathfrak{u})=P(J(u))$, $\operatorname{Det}\left(\mathrm{D}^{2} \phi\right)=\mathrm{P}(\mathrm{J}(\nabla \phi))$.
These relations yield and extend another recent result of S . Müller [5].

3 Proofs

Theorems 1-2 follow immediately from the following lemma and the classical maximal theorem:

Lemma 3: 1) Let $u \in W^{1,4 / 3}\left(\mathbb{R}^{2}\right)^{2}$; then we have for all $t>0, x \in \mathbb{R}^{2}$

$$
\begin{equation*}
\left|h_{t} * J(u)\right| \leq C_{0}\left(f_{B(x, t)}|D u|^{4 / 3} d x\right)^{3 / 2} \tag{6}
\end{equation*}
$$

(2) Let $\phi \in \mathrm{W}^{2,1}\left(\mathbb{R}^{2}\right)$; then we have for all $\mathrm{t}>0, \mathrm{x} \in \mathbb{R}^{2}$

$$
\begin{equation*}
\left|h_{t} * J(\nabla \phi)\right| \leq C_{0}\left(\int_{B(x, t)}\left|D^{2} \phi\right| d x\right)^{2} \tag{7}
\end{equation*}
$$

## Remarks:

1) Here and everywhere below, $C_{0}$ denotes various constants independent of $t, x, u, \phi$.
2) The estimate (7) is still true for $\phi \in \bar{W}^{2,1}$ provided we define $J(\nabla \phi)$ by (4) and the right hand side is replaced by the total mass of the measure $\sum_{i, j}\left|\frac{\partial^{2} \phi}{\partial x_{j} \partial x_{j}}\right|$ on the ball $B(x, t)$.
3) Those estimates allow, in fact, to investigate the borderline cases $\mathrm{p}=\frac{4}{3}$ or $\mathrm{p}=1$.

Proof of Lemma 3: We begin with Part 1). Using (3) and integrating by parts, we find
(8) $\quad h_{t} * J(u)=\int u_{1}(y)\left[\partial_{2} u_{2}\right.$ (y) $\left.\frac{1}{t^{3}} \partial_{1} h\left(\frac{x-y}{t}\right)-\partial_{1} u_{2}(y) \frac{1}{t^{3}} \partial_{2} h\left(\frac{x-y}{t}\right)\right] d y$.

But, clearly, these expressions are left invariant if we subtract constants from $u_{1}$ and $u_{2}$. Therefore, in particular, we have, denoting by $f_{B(x, t)} u=\frac{1}{|B(x, t)|} \int_{B(x, t)} u(y) d y$,

$$
\begin{equation*}
h_{t} * J(u)=\int \frac{1}{t}\left(u_{1}-\int_{B(x, t)} u_{1}\right)\left[\partial_{2} u_{2}(y) \frac{1}{t^{2}} \partial_{1} h\left(\frac{x-y}{t}\right)-\partial_{1} u_{2}(y) \frac{1}{t^{2}} \partial_{2} h\left(\frac{x-y}{t}\right)\right] d y \tag{9}
\end{equation*}
$$

We now apply Hölder's inequality to find

$$
\begin{equation*}
\left|h_{t} * J(u)\right| \leq C_{0}\left(f_{B(x, t)}\left|\frac{1}{t}\left(u_{1}-f_{B(x, t)} u_{1}\right)\right|^{4} d y\right)^{1 / 4}\left(f_{B(x, t)}|D u|^{4 / 3} d y\right)^{3 / 4} \tag{10}
\end{equation*}
$$

And we obtain (6) by recalling the Sobolev-Poincaré's inequality

$$
\begin{equation*}
\left(\left.f_{B(x, t)} \frac{1}{t}\left(u_{1}-f_{B(x, t)} u_{1}\right)\right|^{4} d y\right)^{1 / 4} \leq C_{0}\left(f_{B(x, t)}|D u|^{4 / 3} d y\right)^{3 / 4} \tag{11}
\end{equation*}
$$

We now turn to part 2), which is proven in a similar way using either (4) or (5): we use for instance (4) and obtain as above

$$
\begin{gather*}
\mathrm{h}_{\mathrm{t}} * \mathrm{~J}(\nabla \phi)=-\frac{1}{2} \int\left(\partial_{2} \phi(\mathrm{y})^{2} \frac{1}{\mathrm{t}^{4}} \partial_{1}^{2} \mathrm{~h}\left(\frac{\mathrm{x}-\mathrm{y}}{\mathrm{t}}\right) \mathrm{dy}+\right.  \tag{12}\\
-\frac{1}{2} \int\left(\partial_{1} \phi(\mathrm{y})^{2} \frac{1}{\mathrm{t}^{4}} \partial_{2}^{2} \mathrm{~h}\left(\frac{\mathrm{x}-\mathrm{y}}{\mathrm{t}}\right) \mathrm{dy}+\int \partial_{1} \phi(\mathrm{y}) \partial_{2} \phi(\mathrm{y}) \frac{1}{\mathrm{t}^{4}} \partial_{12}^{2} \mathrm{~h}\left(\frac{\mathrm{x}-\mathrm{y}}{\mathrm{t}}\right) \mathrm{dy}\right.
\end{gather*}
$$

Then we observe that this quantity is left invariant if we add to $\phi$ an arbitrary affine function so that, in other words, we may subtract respectively from $\partial_{1} \phi, \partial_{2} \phi$ the following quantities $\int_{\mathrm{B}(\mathrm{x}, \mathrm{t})} \partial_{1} \phi, \int_{\mathrm{B}(\mathrm{x}, \mathrm{t})} \partial_{2} \phi$.

We then find

$$
\begin{equation*}
\left|h_{t} * J(\nabla \phi)\right| \leq C_{0} f_{B(x, t)}\left|\frac{1}{t}\left\{\nabla \phi-f_{B(x, t)} \nabla \phi\right\}\right|^{2} d y \tag{13}
\end{equation*}
$$

and we conclude again by Sobolev-Poincaré's inequality:

$$
f_{B(x, t)}\left|\frac{1}{t}\{f-\underset{B(x, t)}{f} f\}\right|^{2} d y \leq C_{0}\left(f_{B(x, t)}|D f| d y\right)^{2}, \text { for all } f \in W^{1,1}\left(\mathbb{R}^{2}\right)
$$

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