

ISOPERIMETRIC INEQUALITY FOR SOME
EIGENVALUES OF AN INHOMOGENEOUS,
FREE MEMBRANE

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Abstract

Szegő's inequality concerning the second eigenvalue of a homogeneous, free membrane is extended to the case of an inhomogeneous free membrane. With the help of a variational principle and the conformal mapping technic upper bounds are constructed for the sum $\int U_2 + \int \hat{v}$ where \hat{v} and \hat{v} denote the second and third eigenvalue. These bounds only depend on the total mass of the domain and on a simple expression involving the mass distribution and its logarithm.

(*)

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1. This paper is concerned with the eigenvalue problem

$$(A) \quad \Delta_{\mathbb{R}^2} p + p p c p = 0 \quad \text{in } G \quad \Delta_{\mathbb{R}^2} = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$$

$$p = 0 \quad \text{on } \Gamma,$$

where G is a simply connected domain in the z -plane ($z=x+iy$), Γ is its boundary and n denotes the outer normal. $p(z) > 0$ is the mass distribution. We assume that there exists a countable set of eigenvalues $0 = \lambda_1 < \lambda_2 < \lambda_3 < \dots$

Szegö [5] proved for membranes with constant p the isoperimetric inequality

$$\lambda_n \leq \frac{j_{n-1}^2 + j_{n-2}^2}{4\pi p_1}$$

where $M = \iint_G p \, dx dy$ is the total mass, $J_{n-1}(r)$ denotes the

Bessel function of order $n-1$ and $p_1 = 1.8412\dots$ is the smallest value such that $J_{n-1}'(p_1) = 0$. Equality holds if and only if G

is a circle. In his proof, Szegö uses conformal transplantation.

Weinberger [6] extended this theorem to an N -dimensional not

necessarily simply connected membrane with constant mass distribution. This result can be stated as follows:

For all N-dimensional domains G of given total mass, the N-dimensional sphere yields the maximum of $\int \Delta u$.

The aim of this paper is to generalize the inequality of Szegő to the case of an inhomogeneous membrane.

2. Problem (A) is equivalent to the eigenvalue problem

$$\Delta u + \lambda u = 0 \quad \text{in } G, \quad \frac{\partial u}{\partial \nu} = 0 \quad \text{on } T, \quad \text{where } \Delta = \frac{\Delta^A}{P}$$

Beltrami operator of an abstract surface T given by G and

the line element $ds^2 = p(z) |dz|^2$. Here, $\nu = \nu^i \frac{\partial}{\partial x^i}$ is the unit normal with respect to the metric of \mathbb{R}^n . By the Theorema Egregium it follows that the Gaussian curvature is $K = \frac{G}{p} - (-\Delta \ln p) / 2p$.

We shall assume, that the inequality $(-\Delta \ln p) / 2p < \bar{K}$ holds

in $G \cup T$ and that $2\nu - K M > 0$.

In addition to problem (A) we consider in the w -plane ($w = u + iv$) the auxiliary problem

$$(B) \quad \Delta_w \hat{u} + \lambda \text{jug}(w) \hat{u} = 0 \quad \hat{G} = \{w; |w| < \hat{R}\}$$

$$\frac{\partial \hat{u}}{\partial \rho} = 0 \quad \text{on } \hat{T} = \{w; |w| = \hat{R}\}$$

$$\text{where } g(w) = \begin{cases} \frac{4c^2}{(1 \pm r^2)^2} & \text{if } K_0 = \pm c^{-2} \\ 1 & \text{if } K_0 = 0 \end{cases}$$

[r, θ polar coordinates of the w -plane] and $\hat{n} = \frac{1}{\sqrt{g}} \frac{\partial}{\partial r}$ is the unit normal with respect to the metric $ds^2 = g(w) |dw|^2$. R

is determined such that $\iint_{\hat{G}} g(w) du dv = \iint_{\hat{G}} p(z) dx dy = M$. An

elementary calculation yields

$$\hat{R}^2 = \begin{cases} \frac{M}{47TC^2 + M} & \text{if } K_0 = \pm c^{-2} \\ \frac{M}{\pi} & \text{if } K_0 = 0 \end{cases}$$

Because of our assumptions, we have in the case $K_0 \neq 0$ $R \leq 1$. \hat{G} with the metric ds^2 can be interpreted as a geodesic circle on a surface of constant Gaussian curvature K_0 . The eigenfunctions of problem (B) have the form

$$(2) \quad \hat{\varphi}_k(r, \theta) = R_m(\hat{\mu}_k; r) \begin{cases} \sin m \theta \\ \cos m \theta \end{cases} \quad m = 1, 2, \dots,$$

or

(3) $\phi_k(r, 0) = R_0(l^{\wedge}; r)$

$R_{in}(\hat{L}; r)$ is an eigenfunction of

$$(\hat{r}R) i i \frac{2}{r} R + \epsilon_r g(w) R = 0 \quad \text{in } (0, \hat{R})$$

(4)

$$R'(\hat{R}) = 0, \quad R(0) < 0$$

PROPOSITION. The second eigenvalue $\hat{J}U_2$ $2JL$ problem (B) is degenerated; the corresponding eigenfunctions are $\phi_2(r, \theta) = R_{-1}(\hat{U}_2; r) \sin \theta$ and $\hat{\phi}_-(r, \theta) = R, (\hat{L}; r) \cos \theta$.

Proof. We have to show that the first eigenvalue $\hat{M}_1^{*(1)}$ of (4) corresponding to $m = 1$ is smaller than the second eigenvalue $\hat{M}_2^{(0)}$ of (4) with $m \sim \theta$. If $K_0 = 0$, then $\hat{M}_1^{*(1)} < \hat{M}_2^{(0)}$

and $\hat{W}_2^{(0)} = \left(\frac{2}{R} \right)^2$. Let $K_0 = \pm c^2$ be different from

zero. We introduce the new variable $z = \frac{r^2 - r}{1 + r^2}$ if $K_0 > 0$,

or $z = \frac{r^2}{1 - r}$ if $K_0 < 0$. The interval $[0, \hat{R}]$ is then transformed into $[+1, b+]$ where $b+ = \frac{R}{1 + R^2}$, and the differential equation takes the form

$$(5) \quad \pm \{ (1 - z^2) R' \}' + \frac{m^2}{1 - z^2} \tilde{R} + \hat{\mu} c^2 \tilde{R} = 0$$

with the boundary conditions

$$R(\bar{+} 1) < \infty, \quad R'(b_{\pm}) = 0.$$

If $m = 1$, then the solutions of (5) can be written as

$\tilde{R}(z) = \sqrt{f_{\pm}(1 - z^2)}$ $p'(z)$ [2] where $p(z)$ is a solution of the Legendre equation

$$(6) \quad \pm \{ (1 - z^2) p' \}' + f^{(1)} c^2 p = 0 \quad \text{in } (\bar{+} 1, b_{\pm})$$

with the boundary conditions

$$p(\bar{+} 1) < \infty, \quad p'(b_{\pm}) = \pm (f^{(1)} c b_{\pm}) p(b_{\pm})$$

Assume that $p'(z) < 0$ in $(\bar{+} 1, b_{\pm})$. Because of the boundary condition at b_{\pm} we have $p(b_{\pm}) > 0$, and thus, $p(z) > 0$ in the interval $(\bar{+} 1, b_{\pm})$. According to the definition, $\hat{M}_0^{(0)}$ is the second eigenvalue of $\pm \{ (1 - z^2) p' \}' + j_1^{(0)} c^2 p = 0$ in $(\bar{+} 1, b_{\pm})$ with $p(\bar{+} 1) < \infty$ and $p'(b_{\pm}) = 0$. The corresponding eigenfunction vanishes at some point in $(\bar{+} 1, b_{\pm})$. By the Sturm Comparison Theorem it follows that $\hat{L}_1^{(1)} < \hat{M}_0^{(0)}$.

3. THEOREM, if $(-f \ln p)/2p \leq K_0$ if $G \cup T$ and if $2ir - K_0 M \neq 0$, then the following inequality holds between the eigenvalues of the problems (A) and (B)

$$(7) \quad \frac{1}{\mu_2} + \frac{1}{\mu_3} \geq \frac{1}{\hat{\mu}_2} + \frac{1}{\hat{\mu}_3}$$

Proof. The proof is based on the variational principle

$$(8) \quad \frac{1}{\mu_2} + \frac{1}{\mu_3} = \text{Max}_{v_2, v_3} \{R[v_2, v_3]\}^{-1}$$

$$R[f] = \frac{\int_G f^2 p \, dx dy}{\int_G \text{grad}^2 f \, dx dy}$$

is the Rayleigh quotient- v_2 and v_3 range over all piecewise continuously differentiable functions with

$$D(v_2, v_3) = \int_G \text{grad} v_2 \text{grad} v_3 \, dx dy = 0$$

and

$$\int_G v_i p \, dx dy = 0 \quad i = 2, 3 .$$

Let $z(w)$ be a conformal mapping from G onto \hat{G} , and let $U_2(z) = \hat{c}p_2(w(z))$ and $U_3(z) = \hat{s}_3(w(z))$ [$w(z)$ inverse function of $z(w)$] be the transplanted eigenfunctions. In [5] it is

shown, that there always exists a function $z(w)$ such that $\iint_{U_i} p \, dx dy = 0$ for $i = 2, 3$. Because of the invariance of the Dirichlet integral under conformal mapping, we have

$D_G(U_2, U_3) = D_G(\hat{\phi}_2, \hat{\phi}_3) = 0$. Therefore, U_2 and U_3 are admissible functions for the variational characterization (8). Since $D_{\hat{G}}(U_0) = D_{\hat{G}}(U_2)$, we conclude from (8) that

$$(9) \quad \frac{\iint_{U_i} p \, dx dy}{\int_{\hat{G}} |g|^2 \, r \, dr d\theta} = \frac{\int_{\hat{G}} p(z(w)) |g|^2 \, r \, dr d\theta}{D_{\hat{G}}(\hat{\phi}_2)}$$

We shall use the notations $m(r) = \int_0^{2\pi} p(z(w)) |g|^2 \, r \, dr$,

$\hat{T}_t = \{w; |w| = t\}$, $\hat{G}_t = \{w; |w| < t\}$, $\Gamma_t = z(\hat{T}_t)$ and

$G_t = z(\hat{G}_t)$. $m(t) = \int_0^t m(r) \, dr$ is the total mass of G_t .

Integration by part yields

(1)

In [5] the existence of such a function $z(w)$ has only been proved for $p = 1$, but it is easy to see that the proof remains the same if p is an arbitrary positive and continuous function.

$$\begin{aligned}
 (10) \quad \int_0^{\hat{R}} R_1^2 \frac{d}{dr} m(r) dr &= R_1^2 m(r) \Big|_0^{\hat{R}} - 2 \int_0^{\hat{R}} R_1 \left(\frac{d}{dr} R_1 \right) m(r) dr \\
 &= R_1^2 (R)M - 2 \int_0^{\hat{R}} R_1 \left(\frac{d}{dr} R_1 \right) m(r) dr
 \end{aligned}$$

The next step will be to estimate $m(r)$. Consider the modul $H(T, T_t) = \{D(h)\}^{-1}$. h is the solution of the Dirichlet problem $\Delta h = 0$ in $G \setminus G_t$, $h = 0$ on T and $h = 1$ on T_t . The modul is invariant under conformal transformation, therefore $\mu(\Gamma, r_t) = n(t, f_t) = jft$ in \mathbb{H} .

$$\hat{m}(t) = \int_0^{2\pi} \int_0^t g(w) r dr d\theta = \begin{cases} (4irc^2 t^2 / (1 + t^2)) & \text{if } K = +c^2 \\ \pi t^2 & \text{if } K_0 = 0 \end{cases}$$

denotes the total mass of $G_t = \{w; |w| < t\}$ with respect to the mass distribution of the problem (B). By Corollary 2 [1] we have

$$\begin{aligned}
 (11) \quad \mu(\Gamma, r_t) &= \frac{1}{4\pi} \left\{ \ln \frac{M}{47TC^2 \pm M} - \ln \frac{\hat{m}(t)}{47rc^2 \pm \hat{m}(t)} \right\} \\
 &\leq \frac{1}{4\pi} \left\{ \ln \frac{M}{47TC^2 \pm M} - \ln \frac{m(t)}{47TC^2 \pm m(t)} \right\} \\
 &\text{if } K_0 = \pm c^{-2}
 \end{aligned}$$

or else

$$(12) \quad M(r, r_t) - fcmJ \frac{\Delta}{4\pi} \ln f(\bar{t}) \quad (1)$$

Therefore we obtain the estimation

$$(13) \quad m(t) \leq \hat{m}(t) \quad \text{for all } t.$$

If $r \in [0, \hat{R}]$, then $R - jJr \hat{a} \wedge R_1(r) \wedge 0$. This statement is equivalent with $R_x(z) R[(z) \wedge 0$ in $I = (T1^b + J$ (cf. sec. 1(5)).

It follows immediately from the next result.

LEMMA. Let $f(z)$ be the first eigenfunction of the eigenvalue problem $\{r(z) f'(z)\}' + (A - e(z)) f(z) = 0$ in (a, b) , $f(a) = 0$ and $f(b) = 0$. If $o(z) > 0$ and if $e(z)$ is a non-increasing function, then we have $f(z) f'(z) \wedge 0$ in (a, b) .

This lemma will be proved by contradiction. Since $f(z)$ is the first eigenfunction, it has constant sign in (a, b) . We may assume that $f(z) > 0$. Suppose that $f'(z) < 0$ in some interval. Because of the boundary conditions there exists a

*1' (11) and (12) are generalizations of a theorem by T. Carleman, Math. Z. 1 (1918), pp. 208-212 for the capacity of a condenser. They hold only under the assumptions $(-A \ln p)/2p < K_0$ and $(4\pi - KM) > 0$.

point $x_0 \in (a, b]$ such that $f'(x_0) = 0$ and $f(x_0) > 0$. By multiplying the differential equation with $f(z)$ and integrating,

$$\int_a^{x_0} a(z) f'(z)^2 dz + \int_a^{x_0} (\lambda - e(z)) f^2(z) dz = 0.$$

From this relation it follows that $\lambda_1 > \inf_{z \in (a, x_0)} e(z)$ and because

of the monotonicity of $e(z)$

$$\lambda_1 > e(x_0),$$

At x_0 , $f(z)$ satisfies $f f'(x_0) + (\lambda - e(x_0)) f(x_0) = 0$.

Since $a(x_0)$ and $(\lambda - e(x_0))$ are positive, $f''(z)$ and $f(z)$ must vanish at x_0 . By the uniqueness theorem the only solution for which $f(x_0) = 0$ and $f'(x_0) = 0$ is $f(z) \equiv 0$. But this is no eigenfunction of the eigenvalue problem.

(13) and the monotonicity of $R_1(r)$ together with (10) yield

$$\int_0^{\hat{R}} R_1^2(r) f(r) dr - 2 \int_0^{\hat{R}} R_1^2(r) CD \hat{\phi}_r'(r) dr$$

and hence by (9)

$$\sum_{i=2}^3 \frac{\int_0^{\hat{R}} R_1^2(r) f(r) dr}{Dg(\hat{\phi}_2)} \geq \frac{\int_0^{\hat{R}} R_1^2(r) |\hat{\phi}_r'(r)| dr}{Df(c\hat{\phi}_2)} = \sum_{i=2}^3 \hat{\mu}_i^{-1}.$$

4. The following corollaries are some immediate consequences of the theorem in the previous section.

COROLLARY 1. Consider the eigenvalues $\hat{\mu}_3^{-1}(K_0)$ of problem (B) as a function of the Gaussian curvature K_0 . Suppose that the total mass M is fixed. If $2\pi - K_0 M > 0$, then $\hat{\mu}_3^{-1}(K) + \hat{\mu}_3^{-1}(K_0)$ is a monoton decreasing function of K .

COROLLARY 2. If $2\pi - K_0 M > 0$, then

$$(14) \quad \frac{1}{M_2} + \frac{1}{M_3} \geq \frac{1}{27r}$$

Proof; From the theorem in section 3 and Corollary 1 we have $\frac{1}{\hat{\mu}_2} + \frac{1}{\hat{\mu}_3} \geq \frac{1}{jL(K)} + \frac{1}{\hat{\mu}_3(K)} \geq \frac{1}{\hat{\mu}_0(27r/M)} + \frac{1}{\hat{\mu}_0(27T/M)} = \frac{1}{M}$. $\hat{\mu}_3^{-1}(2\pi/M)$ correspond to the eigenvalues of the half-sphere

with the radius $\sqrt{\frac{M}{27T}}$.

~~REMARK~~ Corollary 2 can also be obtained from the inequality

$$\frac{1}{\lambda_1} + \frac{1}{\mu_2} + \frac{1}{\mu_3} \geq \frac{3M}{\pi} \quad [3] \quad \text{where } \lambda_1 \text{ is the first eigenvalue of}$$

the membrane $\Delta u + \lambda u = 0$ in G , $u = 0$ on T . This result together with the inequality $\hat{\mu}_3^{-1} \geq \hat{\mu}_3^{-1}$ [1] leads to Corollary 2.

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