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## EXPONENTIAL STABILITY FOR

### A NONLINEAR FUNCTIONAL

### DIFFERENTIAL EQUATION

by

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#### 1. Introduction.

This paper concerns the nonlinear functional differential equation,

(E) 
$$y(t) = -\int_{0}^{t} a(t-\tau)g(y(\tau))d\tau + f(t).$$

This equation has been the subject of much study, mainly by Levin and Nohel (see for example [4]) and Hannsgen [2]. The object has been to obtain sufficient conditions for asymptotic stability. The best result in this direction appears in [2] and is contained in the following theorem.

Theorem I. Suppose the following conditions are satisfied:

(A <sub>1</sub> )	$(-1)^{k}a^{(k)}(t) \ge 0, k = 0, 1, 2, a(t) \ne constant, \pm \infty$
(A <sub>2</sub> )	$\xi_g(\xi) > 0$ for $\xi \neq 0$ ; $\int_0^{\xi} g(\xi) d\xi = \infty$ .

$$A_3$$
) f  $\in L_1(0,\infty)$  and either,

(i)  $|f(t)| \leq M$  for all t

<u>or</u> (ii)  $a \in L_1(o, \infty)$  <u>and</u>  $|f(t)| \leq M$  <u>for all</u> t. <u>Then all solutions of</u> (E) <u>tend to zero as</u> t <u>tends to infinity</u>.

Here we consider the rates of decay of solutions and in particular we give conditions for exponential decay. The basic result is as follows.

<u>Theorem II.</u> Suppose  $(A_1) - (A_2)$  hold and in addition that:

(H<sub>1</sub>)  $a(t) - a(\infty) \le Me^{-\alpha t}$ , for some positive constants M and  $\alpha$ , where  $a(\infty)$  denotes the limit of a(t) at infinity,

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- (H<sub>2</sub>)  $|f(t)| \leq Ne^{-\beta t}$ , for some positive constants N and  $\beta$ ,
- $(H_3) \qquad \underline{If} a(\infty) > 0, \underline{then} | f(t) | \le B \quad \underline{for \text{ some positive con-}}$  $\underline{stant} B,$
- and
- (H<sub>4</sub>) g is differentiable in some neighborhood of origin and g'(0) > 0.

<u>Then there exists a constant  $\gamma \leq \alpha$  such that:</u>

- (i) If  $\beta > \gamma$ , every solution of (1) satisfies  $y(t) = O(e^{-\gamma t})$ as  $t \rightarrow \infty$ , for any  $\gamma < \gamma$ ;
- (ii) If  $\beta < \gamma$ , every solution of (1) satisfies  $y(t) = O(e^{-\gamma' t})$ as  $t \rightarrow \infty$ , for any  $\beta' < \beta$ .

Note that the hypotheses of Theorem II imply  $(A_3)$ .

<u>Remarks</u> (1) If  $\beta = \gamma$  one must give more precise information about the behavior of f for large t and we omit the discussion of this case.

(2) As we indicate in section (2), the conditions  $(H_1)$  and  $(H_2)$  are essentially necessary in order to have exponential stability. (3) Both the results and techniques here are similar to those contained in [5]. In [5] more general nonlinearities were admitted but the results were only local. Applied to the present situation they would assert that Theorem (II) holds if |y(0)| is sufficiently small. We emphasize that Theorem (II) is, in contrast, a global result.

(4) The ideas of section (2) are related to those of Halanay [2]. The proofs in [2] are incorrect but have been recently clarified by the authors [6]. 2. The Linearized Equation.

We use a kind of perturbation technique. We rewrite (E) in the form,

(2.1)  $y(t) = - \int_{0}^{t} a(t-\tau)g'(0)y(\tau)d\tau + F(t),$ where, t

(2.2) 
$$F(t) = f(t) - \int_{0}^{t} a(t-\tau) (g(y(\tau)) - g'(0)y(\tau)) d\tau$$

From Theorem (I) we know that y(t) tends to zero. Thus we can consider (E) as a perturbation on the linear equation (2.1). The present section is devoted to a study of (2.1) which is of some interest in itself. We prove the following result. Lemma: Suppose  $(A_1)$ ,  $(H_1)$  and  $(H_4)$  hold. Then any solution of (2.1) can be written in the form:

(2.3) 
$$y(t) = r(t)y(0) + \int_{0}^{t} r(t-\tau) F(\tau) d\tau$$

where the function r(t) \* satisfies,

(2.4)  $|r(t)| \leq \rho e^{-\gamma t}$  for some constants  $\rho$  and  $\gamma$ ,  $0 < \gamma < \alpha$ . Moreover, if  $a(\infty) \neq 0$  then r(t) satisfies, (2.5)  $|\int_{0}^{t} r(\tau) d\tau| \leq \rho e^{-\gamma t}$ ,

where  $\rho'$  is some positive constant.

\* The function r is referred to as the <u>resolvent</u> of equation (2.1).

<u>Proof</u>: We proceed formally by taking the Laplace transform of (2.1). This yields,

(2.6) 
$$s_{y}^{\wedge}(s) - y(0) = -a(s)g'(0)_{y}^{\wedge}(s) + F(s)$$
,

where  $\stackrel{\wedge}{y}$ ,  $\stackrel{\wedge}{a}$  and  $\stackrel{\wedge}{F}$  denote the transforms of y, a and F respectively.

The estimate (H1) yields,

(2.7) 
$$\stackrel{\wedge}{a}(s) = \frac{a(\infty)}{s} + \stackrel{\wedge}{b}(s),$$

where  $\hat{b}$  is analytic in  $\Re e \ s > -\alpha$ . Thus we obtain, from (2.6), (2.8)  $\hat{y}(s) = \hat{r}(s) (y(0) + \hat{F}(s)),$ 

(2.9) 
$$\bigwedge^{\wedge} (s) = \frac{1}{s + a(s)g'(0)} = \frac{1}{s + g'(0)(\frac{a(\infty)}{s} + b(s))}$$

One obtains (2.3), formally, from (2.9) and the convolution theorem; r will be the inverse transform of  $\stackrel{\wedge}{r}$ . To prove estimates (2.4) and (2.5), we must study carefully the function  $\stackrel{\wedge}{r}$ . It is clearly analytic in  $\stackrel{\wedge}{r}$ e s >  $-\alpha$  except at the zeroes of the denominator which are poles. We claim, in fact, that  $\stackrel{\wedge}{r}$  is analytic in  $\stackrel{\wedge}{r}$ e s  $\geq -\gamma$ for some  $\gamma > 0$ . This is proved in [5] and we review, briefly, the argument. It can be shown by a direct calculation that  $(A_1)$  implies the inequality,

It follows that  $s + g'(0)a^{\wedge}(s)$  cannot have zeroes <u>on</u> the imaginary axis. It can be verified, by integration by parts, that

(2.11) 
$$\stackrel{\wedge}{a}(s) = \frac{a(0)}{s} + \frac{a'(0)}{s^2} + o(\frac{1}{s^2}) as |s| \to \infty$$

This equation shows that  $s + g'(0)\hat{a}(s)$  cannot have zeroes of arbitrarily large absolute value. Thus the claim is verified.

Once the properties of  $\stackrel{\wedge}{r}$  are known one can apply the inversion formula for Laplace transforms to obtain r. It is a straightforward matter to verify that (2.3) is indeed a solution of (2.1) and a standard argument shows that this solution is unique. The estimate (2.4) follows from the inversion formula and the fact that  $\stackrel{\wedge}{r}$  is analytic in  $\stackrel{\Re}{\sim}$  s  $\geq -\gamma$ .

The transform of the integral of  $\hat{r}$  is given by  $\hat{r}(s)/s$ . When  $a(\infty) \neq 0$  (2.9) shows that this function has a removable singularity at s = 0 and otherwise shares the analyticity properties of  $\hat{r}$ . Thus (2.5) holds when  $a(\infty) \neq 0$ . This completes the proof of the lemma.

<u>Remark</u>: The estimates  $(H_1)$  and  $(H_2)$  are in some sense necessary for exponential stability. Suppose we have a solution y of the linear equation (2.2) such that  $y(t) = O(e^{-\beta t})$ . Then  $\overset{\wedge}{y}(s)$ is analytic in  $\Re e \ s \ge -\beta$  and hence the right side of (2.8) must be also. But unless both a and F decay exponentially it cannot be that  $\overset{\wedge}{r}$  and  $\overset{\wedge}{F}$  are analytic in such a region. 3. <u>Proof of Theorem II.</u>

Theorem (I) and (H<sub>4</sub>) imply that given any  $\epsilon$  > 0, there exists T > 0 such that, a

$$(3.1) \qquad |g(y(\tau)) - g'(0)y(\tau)| \leq \epsilon |y(\tau)| \quad \text{in } \tau \geq T.$$

We write, for T > 0,

(3.2) 
$$F(t) = F_{m}(t) + G_{m}(t)$$

where,

.

(3.3)  

$$F_{T}(t) = \int_{0}^{T} a(t-\tau) (g(y(\tau)) - g'(0)y(\tau)) d\tau$$

$$G_{T}(t) = \int_{T}^{\infty} a(t-\tau) (g(y(\tau)) - g'(0)y(\tau)) d\tau$$

We observe that (2.3) and (2.4) yield

(3.4) 
$$|\gamma(t)| \leq \gamma(t) + \rho e^{-\gamma t} \int_{0}^{t} e^{\gamma \tau} [F_{T}(\tau) + G_{T}(\tau)] d\tau$$

where

(3.5) 
$$Y(t) = \rho e^{-\gamma t} (|\gamma(0)| + \int_{0}^{t} e^{\gamma t} |f(\tau)| d\tau).$$

. . .

From  $(H_2)$  we have,

Y(t) 
$$\leq Ke^{-\gamma t}$$
 if  $\beta > \gamma$   
3.6)  
Y(t)  $\leq Ke^{-\beta t}$  if  $\beta < \gamma$ .

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Now we split the proof into two cases.

<u>Case</u> (1)  $a(\infty) = 0$ .

In this case (H<sub>1</sub>) shows that  $|F_{T}(\tau)| \leq K e^{-\alpha \tau}$ . Then (3.5) yields, if  $\beta > \gamma$ ,

(3.7) 
$$|_{Y}(t)| \leq K''e^{-\gamma t} + \rho e^{-\gamma t} \int_{O}^{t} e^{\gamma \tau} G_{T}(\tau) d\tau$$

With  $\varepsilon>0$  given, choose T so that (3.1) holds. Then for t>T we obtain,

+

$$(3.8) | | y(t) | \leq K'' e^{-\gamma t} + \rho \in Me^{-\gamma t} \int_{O}^{t} e^{\gamma \tau} (\int_{T}^{\tau} e^{-\alpha (\tau - \mu)} | y(\mu) | d\mu) d\tau$$

$$\leq K'' e^{-\gamma t} + \rho \in Me^{-\gamma t} \int_{O}^{t} | y(\mu) | (\int_{\mu}^{t} e^{\gamma \tau} e^{-\alpha (\tau - \mu)} d\tau) d\mu$$

$$\leq K'' e^{-\gamma t} + \frac{\rho \in M}{\alpha - \gamma} e^{-\gamma t} \int_{O}^{t} | y(\mu) | e^{\gamma \mu} d\mu.$$

Set 
$$z(t) = e^{\gamma t} |y(t)|$$
 so that (3.8) yields,  
(3.9)  $z(t) \leq K'' + \nu \int_{0}^{t} z(\mu) d\mu$   $\nu = \frac{\rho \in M}{\alpha - \gamma}$ 

Since  $\ \varepsilon$  is arbitrary we deduce that

$$z(t) = O(e^{mt})$$

for arbitrary m > 0 and hence

$$y(t) = O(e^{-\gamma' t})$$
 for any  $\gamma' < \gamma$ .

If  $\beta < \gamma$ , similar calculations yield

$$y(t) = O(e^{-\beta't})$$
 for any  $\beta' < \beta$ .

<u>Case</u> (2)  $a(\infty) \neq 0$ .

The proof proceeds in almost the same way. The only difference is that we no longer have the estimate  $|F_{T}(\tau)| \leq K'e^{-\alpha \tau}$ . Thus we need a preliminary step. We have, by (H<sub>1</sub>),

$$F(\tau) = a(\infty) \int_{0}^{\tau} (g(y(\mu)) - g(0) y(\mu)) d\mu + \widetilde{F}(\tau)$$

where  $\widetilde{F}(\tau)$  is just like the term treated in case (1). For the first term we have,

$$a(\omega) \int_{O}^{t} r(t-\tau) \int_{O}^{\tau} (g(y(\mu)) - g'(O)) y(\mu) d\mu d\tau = a(\omega) \int_{O}^{t} r(\tau) \int_{O}^{t-\tau} (g(y(\mu)) - g'(O) y(\mu)) d\mu d\tau$$

$$= a(\omega) \left\{ \left( \int_{O}^{\xi} r(\xi) d\xi \right) \int_{O}^{t-\tau} (g(y(\mu)) - g'(O) y(\mu)) d\mu \right\}_{O}^{T}$$

$$+ \int_{O}^{t} (\int_{O}^{\tau} r(\xi) d\xi) (g(y(t-\tau)) - g'(O) y(t-\tau)) d\tau \right\}$$

$$= a(\omega) \int_{O}^{t} (\int_{O}^{t-\tau} r(\xi) d\xi) (g(y(\tau)) - g'(O) y(\tau)) d\tau .$$

By (2.5) of the lemma we have,  

$$\begin{aligned} t-\tau \\ \left| \int_{0} (r(\xi) d\xi) \right| &\leq \rho \cdot e^{-\gamma (t-\tau)} \end{aligned}$$

and now we can proceed as in the preceding case.

<u>Remark</u>. The assumption  $(H_4)$  was dictated the method of proof. This linearization process is used in order to obtain a global result. Without  $(H_4)$ , one can still obtain a local result, see [5]. On the other hand, without exponential decay of the resolvent r(t), linearization still yields a local result for asymptotic stability, in fact for more general equations, see [7].

#### 4. An Application to Control Theory.

Consider the autonomous system of n + 1 ordinary differential equations,

(4.1) 
$$\begin{cases} x = Ax - \alpha g(y), \\ \cdot \\ y = \beta x - \gamma g(y), \end{cases}$$

where  $\mathbf{x}, \boldsymbol{\alpha}$  are n-dimensional column vectors, A is an  $n \times n$  matrix,  $\beta$  is an n-dimensional row vector, and  $\mathbf{y}, \mathbf{y}$  are scalars. We solve the first n equations of (4.1) by the variation of constants formula and substitute into the last equation of (4.1). This yields,

(4.2) 
$$y(t) = -\gamma g(y(t)) - \int_{0}^{t} \beta e^{A(t-\tau)} \alpha g(y(\tau)) d\tau + \beta e^{At} x(0).$$

The problem of absolute stability in the theory of nonlinear controls [3] is to determine conditions on A,  $\alpha$ ,  $\beta$ ,  $\gamma$  and g so that all solutions y(t) of (4.2) tend to zero as t tends to infinity. In the special case of direct control, that is when  $\gamma = 0$ , equation (4.2) is of the form (E) and the present result concerning exponential stability becomes applicable. It is usually assumed that g satisfies (A<sub>2</sub>) and A is a stable matrix. Suppose further that  $\beta = \alpha^{T}$ , the transpose of  $\alpha$ . Finally assume that g satisfies (H<sub>5</sub>) and that A is symmetric. Then one can show that  $a(t-\tau) = \alpha^{T}e^{A(t-\tau)}\alpha$  satisfies (H<sub>1</sub>) and (H<sub>2</sub>). Hence it follows from our theorem that the control system (4.1) is not only absolutely stable but exponentially stable.

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