

NOTICE WARNING CONCERNING COPYRIGHT RESTRICTIONS:
The copyright law of the United States (title 17, U.S. Code) governs the making of photocopies or other reproductions of copyrighted material. Any copying of this document without permission of its author may be prohibited by law.

NONLINEAR EIGENVALUE PROBLEMS AND
CRITICAL POINTS OF FUNCTIONS

James M. Greenberg

Report 68-39

University Libraries
Carnegie Mellon University
Pittsburgh PA 15213-3890

1. Introduction

We here study the nonlinear eigenvalue problem

$$A \underline{x} + \underline{F}(\underline{x}) = \lambda \underline{x} \quad (1.1)$$

where $A : \mathbb{R}^n \rightarrow \mathbb{R}^n$ is self adjoint and linear and $\underline{F}(\cdot)$ is the gradient of a potential Ψ ; i.e.

$$\underline{F}(\underline{x}) = \nabla_{\underline{x}} \Psi(\underline{x}) . \quad (1.2)$$

It is well known that the nontrivial solutions of (1.1) of fixed amplitude r (i.e. $\underline{x} \cdot \underline{x} = r^2$) are the critical points of $\phi(\underline{x}) \equiv \frac{1}{2} A \underline{x} \cdot \underline{x} + \Psi(\underline{x})$ on $\underline{x} \cdot \underline{x} = r^2$. Moreover, if \underline{y} is such a critical point, then the eigenvalue $\lambda(\underline{y})$ is given by

$$\lambda(\underline{y}) = \frac{(A\underline{y} + \nabla_{\underline{x}} \Psi(\underline{y})) \cdot \underline{y}}{r^2} . \quad (1.3)$$

It is no loss in generality to assume that

$$\frac{1}{2} A \underline{x} \cdot \underline{x} = \sum_{i=1}^n \lambda_i x_i^2 ; \quad (1.4)$$

in this case ϕ takes the form

$$\phi(\underline{x}) \equiv \sum_{i=1}^n \lambda_i x_i^2 + \Psi(\underline{x}) . \quad (1.5)$$

Our interest is in showing that if appropriate conditions are met, then the nontrivial solutions of (1.1) (eigenvectors of $A + F$) may be parameterized smoothly by r .

We also obtain results about the maximal extension of a given branch of eigenvectors.

MAR 21 '69

HUNT LIBRARY
CARNEGIE-MELLON UNIVERSITY

2. Statement of Results

Let ϕ be given by (1.4) and assume that the numbers $\{\lambda_j, j = 1, 2, \dots, n\}$ are distinct and indexed in decreasing order ($\lambda_j > \lambda_{j+1}$). In addition assume that the map $\underline{x} \rightarrow \Psi(\underline{x}) : \mathbb{R}^n \rightarrow \mathbb{R}$ is smooth (C^3 will suffice) and satisfies

$$\left(\sum_{|\alpha|=j} |D^\alpha \Psi(\underline{x})|^2 \right)^{1/2} \leq K \|\underline{x}\|^{3-j}, \quad j = 0, 1, 2. \quad (2.1)$$

In (2.1) D^α stands for any derivative of order j and $\|\cdot\|$ for the Euclidean norm.

The assumption that the λ_j 's are distinct implies that the vectors

$$\pm r \underline{e}_j, \quad \underline{e}_j = (\overbrace{0, \dots, 0}^{j-1}, \overbrace{1, 0, \dots, 0}^j, \overbrace{0, \dots, 0}^{n-j}), \quad j = 1, 2, \dots, n,$$

are the unique critical points of $\phi_0 \equiv \sum_{i=1}^n \lambda_i x_i^2$ on $\underline{x} \cdot \underline{x} = r^2$.

For each $0 < \epsilon < 1$ we let

$$\eta_j^{+(-)}(1, \epsilon) = \left\{ \underline{v} \mid \begin{matrix} v_j = + \\ (-) \end{matrix} \sqrt{1 - \sum_{\substack{k=1 \\ k \neq j}}^n v_k^2}, \sum_{\substack{k=1 \\ k \neq j}}^n v_k^2 \leq \epsilon^2 \right\}. \quad (2.2) \quad *$$

For our purposes we will want two numbers $0 < \epsilon_1 < \epsilon_2 < 1$ such that

$$* \text{ or equivalently } \eta_j^{+(-)}(1, \epsilon) \equiv \left\{ \underline{v} \mid \|\underline{v}\|^2 = 1, \left\| \underline{v} \begin{matrix} - \\ (+) \end{matrix} \underline{e}_j \right\|^2 \leq \epsilon^2 \right\}$$

the neighborhoods $\{\eta_j^{+(-)}(1, \epsilon_1), j = 1, 2, \dots, n\}$ are disjoint on the unit sphere while the neighborhoods $\{\eta_j^{+(-)}(1, \epsilon_2), j = 1, 2, \dots, n\}$ cover the unit sphere.

Theorem 1 (Local Existence and Uniqueness Theorem)

There is an $r_0 > 0$ such that for any $r \in (0, r_0]$ the function ϕ has exactly $2n$ critical points on the sphere $\underline{x} \cdot \underline{x} = r^2$. These points may be labeled in pairs $(\underline{x}_{\sim j}^+(r), \underline{x}_{\sim j}^-(r)) \equiv r(\underline{v}_{\sim j}^+(r), \underline{v}_{\sim j}^-(r))$ according to the scheme

$$\underline{v}_{\sim j}^+(r) \in \eta_j^+(1, \epsilon_1) \text{ and } \underline{v}_{\sim j}^-(r) \in \eta_j^-(1, \epsilon_1). \quad (2.3)$$

The functions $r \rightarrow \underline{v}_{\sim j}^+(r)$ (respectively $r \rightarrow \underline{v}_{\sim j}^-(r)$) are $C^1(0 < r \leq r_0)$ and satisfy

$$\lim_{r \rightarrow 0} \underline{v}_{\sim j}^+(r) = \underline{e}_j \text{ and } \lim_{r \rightarrow 0} \underline{v}_{\sim j}^-(r) = -\underline{e}_j, \quad j = 1, 2, \dots, n. \quad (2.4)$$

In order to state the global existence theorem it is necessary to introduce some additional notation. For each $\underline{v} \ni \|\underline{v}\|^2 = 1$ we let $V(\underline{v})$ be the $n-1$ dimensional vector space

$$V(\underline{v}) \equiv \{\underline{u} \in \mathbb{R}^n \mid \underline{u} \cdot \underline{v} = 0\} \quad (2.5)$$

For each $r > 0$ and $\underline{v} \ni \|\underline{v}\|^2 = 1$ we define the symmetric bilinear form $B(r\underline{v}; \cdot, \cdot) : V(\underline{v}) \times V(\underline{v}) \rightarrow \mathbb{R}$ by

$$B(r\underline{v}; \underline{u}, \underline{w}) \equiv \frac{1}{r^2} \left. \frac{\partial^2}{\partial s \partial t} \phi(r(\underline{v} + s\underline{u} + t\underline{w})) \right|_{s=t=0} - \frac{1}{r} (\nabla_{\underline{x}} \phi(r\underline{v}) \cdot \underline{v}) (\underline{u} \cdot \underline{w}). \quad (2.6)$$

$\mathfrak{B}(rv): V(\underline{v}) \rightarrow V(\underline{v})$ is the linear operator associated with the bilinear form $B(rv; \cdot, \cdot)$.*

Theorem 2 (Global Existence of a Given Branch of Critical Points)

For each j it is possible to extend the function $r \rightarrow \underline{v}_j^+(r)$ from $[0, r_0]$ to some maximal interval $[0, R_j^+)$ in such a way that the function $\underline{x}_j^+(r) \equiv r \underline{v}_j^+(r)$ is a critical point of ϕ (or $\underline{x} \cdot \underline{x} = r^2$). The function $\underline{v}_j^+(\cdot)$ is extended as the unique solution of the initial value problem:

$$\begin{aligned} \mathfrak{B}(rv) \dot{\underline{v}}(r) &\equiv \underline{F}(r, \underline{v}), \quad r > r_0 \\ \underline{v}(r_0) &\equiv \underline{v}_j^+(r_0) \end{aligned} \tag{2.7}$$

The initial data $r_0 \underline{v}_j^+(r_0)$ is the unique critical point of ϕ on $\underline{x} \cdot \underline{x} = r_0^2$ such that $\underline{v}_j^+(r_0) \in \eta_j^+(1, \epsilon)$, and $\underline{F}(r, \underline{v}) \in V(\underline{v})$ is defined by

$$\underline{F}(r, \underline{v}) \equiv \frac{\partial}{\partial r} \left\{ \frac{1}{r} \left[\nabla_{\underline{x}} \Psi(r, \underline{v}) - (\nabla_{\underline{x}} \Psi(r, \underline{v}) \cdot \underline{v}) \underline{v} \right] \right\}.$$

The number R_j^+ is characterized as the first $r > 0$ such that the quadratic form $B(rv_j^+(r); \underline{u}, \underline{u})$ has zero as a critical value on the unit sphere $V(\underline{v}_j^+(r))$. For all $r < R_j^+$ the quadratic has $j-1$ positive and $n-j$ negative critical values on the unit sphere $V(\underline{v}_j^+(r))$.

The following example shows that Theorems 1 and 2 are the best that may be expected.

* For any $\underline{u} \in V(\underline{v})$ $B(rv)\underline{u} = [\nabla_{\underline{x}} \nabla_{\underline{x}} \phi(rv)]\underline{u} - (\underline{u} \cdot [\nabla_{\underline{x}} \nabla_{\underline{x}} \phi(rv)] \underline{v})\underline{v} - \frac{1}{r} (\nabla_{\underline{x}} \phi(rv) \cdot \underline{v})\underline{u}$.

Let

$$\phi_0 = x^2 + 2y^2 ,$$

$$\Psi = \frac{2}{3} x^3 - (x^2 + y^2)(x^2 + 2y^2) ,$$

and

$$\phi = \phi_0 + \Psi .$$

If we introduce the polar coordinates

$$x = r \cos \theta \quad \text{and} \quad y = r \sin \theta ,$$

then $\eta(r, \theta) \equiv \frac{\phi}{r^2}(r \cos \theta, r \sin \theta)$ takes the form

$$\eta(r, \theta) \equiv (1-r^2) (\cos^2 \theta + 2 \sin^2 \theta) + \frac{2}{3} r \cos^3 \theta .$$

A simple computation shows that the critical points of η on $x^2 + y^2 = r^2$ are those numbers $\theta \in [0, 2\pi)$ which satisfy

$$\frac{\partial \eta}{\partial \theta}(r, \theta) = \sin 2\theta (1 - r^2 - r \cos \theta) = 0 .$$

It is clear that for all $r > 0$ the numbers $\theta = 0, \pi/2,$

$\pi,$ and $3\pi/2$ are critical points of η . For

$\frac{\sqrt{5}-1}{2} < r < \frac{\sqrt{5}+1}{2}$ there are two additional critical points

$\theta_1 \in (0, \pi)$ and $\theta_2 \in (\pi, 2\pi)$ which satisfy

$$\cos \theta_i = \frac{1 - r^2}{r} , \quad \frac{\sqrt{5}-1}{2} < r < \frac{\sqrt{5}+1}{2} , \quad i = 1, 2 .$$

Finally, for $r > \frac{\sqrt{5}+1}{2}$ the numbers $\theta = 0, \frac{\pi}{2}, \pi,$ and $\frac{3\pi}{2}$ are again the

only critical points.

We now analyze these critical points in more detail. For $0 < r < \frac{\sqrt{5}-1}{2}$ the points $\theta = 0$ and π correspond to relative minima of η while $\theta = \pi/2$ and $3\pi/2$ correspond to relative maxima. At $r = \frac{\sqrt{5}-1}{2}$ (the point where $\frac{\partial^2 \eta}{\partial \theta^2} \Big|_{\theta=0} = 0$) the

point $\theta = 0$ becomes an inflection point, and for all $r > \frac{\sqrt{5}+1}{2}$ is a maximum.

For $\frac{\sqrt{5}-1}{2} < r < 1$ the points $\theta_1 \in (0, \pi/2)$ and $\theta_2 \in (3\pi/2, 2\pi)$ correspond to relative minima of η and the character of the points $\theta = \frac{\pi}{2}, \pi,$ and $3\pi/2$ is as before. At $r = 1$ (the point where

$\frac{\partial^2 \eta}{\partial \theta^2} \Big|_{\theta = \pi/2} = \frac{\partial^2 \eta}{\partial \theta^2} \Big|_{\theta = 3\pi/2} = 0$) the point θ_1 coalesces with $\pi/2$.

and θ_2 with $3\pi/2$. For $1 < r < \frac{\sqrt{5}+1}{2}$ $\theta = \frac{\pi}{r}$ and $\frac{3\pi}{2}$ become relative minima, the points $\theta_1 \in (\pi, 3\pi/2)$ relative maxima. $\theta = \pi$ is still a relative minimum. At $r = \frac{\sqrt{5}+1}{2}$ (the place where

$\frac{\partial^2 \eta}{\partial \theta^2} \Big|_{\theta = \pi} = 0$) $\theta_1 = \theta_2 = \pi$ is an inflection point of η . For

$r > \frac{\sqrt{5}+1}{2}$ the critical points $\theta = 0$ and π correspond to relative maxima of η while $\theta = \pi/2$ and $3\pi/2$ correspond to relative minima.

3. Proofs

To establish Theorem 1 we look at ϕ in neighborhoods of the critical points of $\phi_0 \equiv \sum_{i=1}^n \lambda_i x_i^2$ on $\tilde{x} \cdot \tilde{x} = r^2$. Clearly, it suffices to look at ϕ in a neighborhoods of re_1 .

We now let $0 < \epsilon_1 < \epsilon_2 < 1$ be two numbers such that the neighborhoods $\eta_j^{+(-)}(1, \epsilon_1)$, $j = 1, 2, \dots, n$ are disjoint on $\tilde{v} \cdot \tilde{v} = 1$ while the neighborhoods $\eta_j^{+(-)}(1, \epsilon_2)$, $j = 1, 2, \dots, n$ cover $\tilde{v} \cdot \tilde{v} = 1$. We set $\tilde{x} = r\tilde{v}$, $\tilde{v} \equiv (v_1, v_2, \dots, v_n)$ and introduce local coordinates:

$$v_1 = \sqrt{1 - \sum_{k=2}^n v_k^2}$$

where $\sum_{k=2}^n v_k^2 \leq \epsilon_2^2$. (3.1)

The function ϕ becomes $r^2 \hat{\phi}(r, \tilde{v})$ where

$$\begin{aligned} \hat{\phi}(v_2, v_3, \dots, v_n; r) &\equiv [\lambda_1 + \sum_{k=2}^n \mu_k v_k^2] \\ &+ \frac{1}{r^2} \Psi(r \sqrt{1 - \sum_{k=2}^n v_k^2}, rv_2, \dots, rv_n), \end{aligned} \quad (3.2)$$

and $\mu_k^1 = \lambda_k - \lambda_1$.

To obtain Theorem 1 it suffices to show that for r sufficiently small (\leq some r_0)

(A) There exists a unique $n-1$ tuple $(v_2^+, v_3^+, \dots, v_n^+)$ with

$$\sum_{k=2}^n v_k^2 \leq \epsilon_1^2 \text{ such that}$$

$$\begin{aligned} & \frac{\partial \hat{\phi}}{\partial v_i} (v_2, v_3, \dots, v_n; r) \\ &= \left\{ 2\mu_i^1 - \frac{\psi_{x_1} \left(r \sqrt{1 - \sum_{k=2}^n v_k^2}, rv_2, \dots, rv_n \right)}{r \sqrt{1 - \sum_{k=2}^n v_k^2}} \right\} v_i \\ &+ \frac{1}{r} \psi_{x_i} \left(r \sqrt{1 - \sum_{k=2}^n v_k^2}, rv_2, \dots, rv_n \right) \equiv 0, \quad i = 2, 3, \dots, n. \quad (3.3) \end{aligned}$$

(B) The $(n-1)$ tuple $(v_2^+, v_3^+, \dots, v_n^+)$ satisfying (3.3) is the only solution in the larger sphere $\sum_{k=1}^n v_k^2 \leq \epsilon_2^2$; and

(C) the map $r \rightarrow (v_2^+(r), v_3^+(r), \dots, v_n^+(r))$ is $C^1(0 < r \leq r_0)$ and satisfies

$$\lim_{r \rightarrow 0} (v_2^+(r), v_3^+(r), \dots, v_n^+(r)) = (0, 0, \dots, 0). \quad (3.4)$$

The growth condition on ψ implies that

$$\left| \frac{\frac{1}{r} \psi_{x_1} \left(r \sqrt{1 - \sum_{k=2}^n v_k^2}, rv_2, \dots, rv_n \right)}{\sqrt{1 - \sum_{k=2}^n v_k^2}} \right| \leq \frac{K r}{\sqrt{1 - \epsilon_2^2}} \quad (3.5)$$

and

$$\left(\sum_{i=2}^n \left| \frac{\psi_{x_i} \left(r \sqrt{1 - \sum_{k=2}^n v_k^2}, rv_2, \dots, rv_n \right)}{r} \right|^2 \right)^{1/2} \leq K r \quad (3.6)$$

for all (v_2, v_3, \dots, v_n) satisfying

$$\sum_{k=2}^n v_k^2 \leq \epsilon_2^2 \quad (3.7)$$

Equation (3.5) implies that if

$$r \leq r_1 \equiv \frac{\min_i |\mu_i^1| \sqrt{1 - \epsilon_2^2}}{K} = \frac{|\mu_2^1| \sqrt{1 - \epsilon_2^2}}{K}, \quad (3.8)$$

and if (3.7) holds, then the operator

$$P \equiv \text{diag} (p_1, p_2, \dots, p_n)$$

with

$$p_i = 2\mu_i^1 - \frac{\psi_{x_1} \left(r \sqrt{1 - \sum_{k=2}^n v_k^2}, rv_2, \dots, rv_n \right)}{r}$$

is invertible and

$$\|P^{-1}\| \leq \frac{1}{|\mu_2^1|} \quad (3.9)$$

It now follows that solving (3.3) is equivalent to solving

$$v_i \equiv T_i(v_2, v_3, \dots, v_n; r), \quad i = 2, 3, \dots, n \quad (3.10)$$

where

$$T_i(v_2, v_3, \dots, v_n; r) \equiv \frac{1}{2\mu_i^1 - \frac{\psi_{x_1}}{r}} \psi_{x_1}(r \sqrt{1 - \sum_{k=2}^n v_k^2}, \quad rv_2, \dots, rv_n) \quad (3.11)$$

We now observe that for $r \leq r_1$ and (v_2, v_3, \dots, v_n) satisfying (3.7)

$$\left(\sum_{k=2}^n T_k^2 \right)^{1/2} \leq \frac{K r}{|\mu_2^1|} \quad (3.12)$$

Equation (3.12) implies that for

$$r \leq \min\left(r_1, \frac{|\mu_2^1| \epsilon_1}{K}\right) \equiv r_2 \quad (3.13)$$

$$T(\cdot, r) : \sum_{k=2}^n v_k^2 \leq \epsilon_2^2 \rightarrow \sum_{k=2}^n v_k^2 \leq \epsilon_1^2 ; \quad (3.14)$$

hence for $r \leq r_2$ any fixed point of (3.10) in $\sum_{k=2}^n v_k^2 \leq \epsilon_2^2$

must be in $\sum_{k=2}^n v_k^2 \leq \epsilon_1^2$. The smoothness of ψ implies that

$T(\cdot, r)$ is C^1 in (v_2, v_3, \dots, v_n) and hence Brouwer's Theorem

guarantees (for all $r \leq r_2$) the existence of at least one solution of (3.10) (and hence (3.3)) in $\sum_{k=2}^n v_k^2 \leq \epsilon_2^2$.

To establish uniqueness, it suffices to show that for some $r_0 \leq r_2$ and all r in $[0, r_0]$ the maps $T(\cdot, r)$ are contractions on

$\sum_{k=2}^n v_k^2 \leq \epsilon_2^2$. This computation follows from (2.1).

The smoothness of $r \rightarrow \underline{v}_1(r) \equiv (\sqrt{1 - \sum_{k=2}^n v_k^2(r)}, v_2(r), \dots, v_n(r))$

follows from the smoothness of $\underline{x} \rightarrow \underline{\psi}(\underline{x})$. We find that for $0 < r \leq r_0$

$\dot{v}_i(r) \equiv \frac{d}{dr} v_i(r)$ exists, is continuous, and satisfies

$$\sum_{j=2}^n B_{ij}(v_2, v_3, \dots, v_n, r) \dot{v}_j(r) = F_i(v_2, v_3, \dots, v_n; r), \quad i = 2, 3, \dots, n, \quad (3.15)$$

where

$$B_{ij}(v_2, v_3, \dots, v_n; r)$$

$$= \left\{ 2\mu_i^1 - \frac{\psi_{x_1}}{r \sqrt{1 - \sum_{k=2}^n v_k^2}} \right\} (r \sqrt{1 - \sum_{k=2}^n v_k^2}, rv_2, \dots, rv_n) \delta_{ij}$$

$$- \frac{\{\psi_{x_1 x_j} v_i + \psi_{x_1 x_i} v_j\}}{\sqrt{1 - \sum_{k=2}^n v_k^2}} (r \sqrt{1 - \sum_{k=2}^n v_k^2}, rv_2, \dots, rv_n)$$

$$\begin{aligned}
 & + \frac{v_i v_j}{(1 - \sum_{k=2}^n v_k^2)^{3/2}} \left\{ \sqrt{1 - \sum_{k=2}^n v_k^2} \Psi_{x_1 x_1} - \frac{\Psi_{x_1}}{r} \right\} (r \sqrt{1 - \sum_{k=2}^n v_k^2}, rv_2, \dots, rv_n) \\
 & + \Psi_{x_1 x_j} (r \sqrt{1 - \sum_{k=2}^n v_k^2}, rv_2, \dots, rv_n); \quad i, j = 2, 3, \dots, n, \quad (3.16)
 \end{aligned}$$

and

$$F_i(v_2, v_3, \dots, v_n; r) \equiv$$

$$\frac{1}{r^2} \left\{ \Psi_{x_i} - \frac{\Psi_{x_1} v_i}{\sqrt{1 - \sum_{k=2}^n v_k^2}} \right\} (r \sqrt{1 - \sum_{k=2}^n v_k^2}, rv_2, \dots, rv_n)$$

$$\frac{v_i}{r} \left\{ \Psi_{x_1 x_1} + \frac{\sum_{j=2}^n \Psi_{x_1 x_j} v_j}{\sqrt{1 - \sum_{k=2}^n v_k^2}} \right\} (r \sqrt{1 - \sum_{k=2}^n v_k^2}, rv_2, \dots, rv_n)$$

$$+ \frac{1}{r} \left\{ \sqrt{1 - \sum_{k=2}^n v_k^2} \Psi_{x_1 x_i} + \sum_{j=2}^n \Psi_{x_i x_j} v_j \right\} (r \sqrt{1 - \sum_{k=2}^n v_k^2}, rv_2, \dots, rv_n),$$

$$i = 2, 3, \dots, n. \quad (3.17)$$

To obtain the limiting relation (3.4) we simply make use of the estimate (3.12).

To establish theorem 2 it again suffices to work with a particular branch of critical points of ϕ . We shall extend the branch

$x_1^+(r) \equiv r v_1^+(r)$, $r \in [0, r_0]$. It is clear that if we extend

$v_1^+(r)$ as a solution the initial value problem (2.7), then

$x_1^+(r) \equiv r v_1^+(r)$ will be a critical point of ϕ on $\underline{x} \cdot \underline{x} = r^2$

It is also clear that the initial value problem (or any of its representations) has a (have) unique solution(s) provided the operator $\mathfrak{B}(r v_1^+(r)) : V(v_1^+(r)) \rightarrow V(v_1^+(r))$ is invertible.

The condition for the lack of invertibility \mathfrak{B} along $r v_1^+(r)$

is simply that the quadratic $B(r v_1^+(r) ; \underline{u}, \underline{u})$ have 0 as a

critical value on the unit sphere $V(v_1^+(r))$. That $B(r v_1^+(r) ; \underline{u}, \underline{u})$

has no positive and $n-1$ negative critical values for $r < R_1$

follows from the fact that $\mathfrak{B}(r v_1^+(r))$ is symmetric and invertible for

$r < R_1$ and the fact that $\mathfrak{B}_{0,1} \equiv \lim_{r \rightarrow 0^+} \mathfrak{B}(r v_1^+(r)) = \text{diag} (\mu_2^1, \mu_3^1, \dots, \mu_n^1)$

maps $V(\underline{e}_1) \rightarrow V(\underline{e}_1)$ and has no positive and $n-1$ negative eigenvalues.